

ENERGY DECAY FOR THE WAVE EQUATIONS IN EXTERIOR DOMAINS

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ABSTRACT. We study the decay estimates of the energy for the wave equation in an exterior domain with a localized dissipation. The dissipative term consists of the following two parts: The first part may be nonlinear and localized in suitable bounded area, while the second part is linear in the outside of a big ball. So we may call such a dissipation as 'half-linear' dissipation.

1. INTRODUCTION

In this paper we consider the initial-boundary value problem for the wave equations with a dissipation ;

$$(1.1) \quad u_{tt} - \Delta u + \rho(x, u_t) = 0 \quad \text{in } \Omega \times (0, \infty)$$

$$(1.2) \quad u(x, 0) = u_0, \quad u_t(x, 0) = u_1 \quad \text{in } \Omega$$

$$(1.3) \quad u(x, t) = 0 \quad \text{on } \partial\Omega \times (0, \infty),$$

where Ω is an exterior domain in \mathbb{R}^N such that $V \equiv \mathbb{R}^N \setminus \Omega$ is compact, the boundary $\partial\Omega$ is smooth, say C^3 class, and $\rho(x, v)$ is a function like $\rho(x, v) = a(x)g(v)$ with $a(x) \geq 0$ and $g'(v) \geq 0$. When $\Omega = \mathbb{R}^N$, Cauchy problem, the boundary condition (1.3) should be dropped.

To state our precise assumptions on $\rho(x, v)$, we introduce a part of the boundary $\partial\Omega$ as follows (see Lions [2]):

$$\Gamma(x_0) = \{x \in \partial\Omega \mid (x - x_0) \cdot \nu(x) > 0\},$$

where $x_0 \in \mathbb{R}^N$ is arbitrarily fixed and $\nu(x)$ denotes the outward unit normal vector at $x \in \partial\Omega$.

We set $B_r := \{x \in \mathbb{R}^N \mid |x| < r\}$ and $\Omega_r := \Omega \cap B_r$.

Let $a(x)$ be a nonnegative bounded function on Ω , satisfying:

(a) There exist $x_0 \in \mathbb{R}^N$ and a relatively open set ω in $\overline{\Gamma(x_0)}$ such that

$$(1.4) \quad \overline{\Gamma(x_0)} \subset \omega \text{ and } a(x) \geq \varepsilon_0 > 0 \text{ a.e. in } \omega;$$

(b) There exists $R > 0$ such that

$$(1.5) \quad a(x) \geq \varepsilon_0 > 0 \text{ a.e. for } |x| \geq L.$$

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With $a(x)$ above we now make the hypotheses on $\rho(x, v)$.

Hyp.A $\rho(x, v)$ is a monotone increasing and differentiable function in $v (\neq 0)$ and satisfies the following conditions:

$$(1.6) \quad \rho(x, v) = a(x)v \text{ if } (x, v) \in B_R^C \times \mathbb{R},$$

for some $R > 0$. If $(x, v) \in \Omega_R \times \mathbb{R}$ and $|v| \leq 1$, then

$$(1.7) \quad k_0 a(x)|v|^{p+2} \leq \rho(x, v)v \leq k_1 a(x)\{|v|^{p+2} + |v|^2\},$$

and if $(x, v) \in \Omega_R \times \mathbb{R}$ and $|v| \geq 1$, then

$$(1.8) \quad k_0 a(x)|v|^{q+2} \leq \rho(x, v)v \leq k_1 a(x)\{|v|^{q+2} + |v|^2\},$$

where $k_0, k_1 > 0$.

We note that our dissipation may be localized near a part of the boundary $\Gamma(x_0)$ and near infinity and may vanish in a large area. The main purpose of this paper is to derive a decay estimate of the energy for the problem (1.1)-(1.3).

When $\rho(x, u_t) = a(x)u_t$, linear in the whole domain, decay properties of the solutions of (1.1)-(1.3) have been investigated by one of the present authors in [7] and it is an interesting problem to generalize them to the case of nonlinear dissipation like $\rho(x, u_t) = a(x)|u_t|^r u_t$. The problem with nonlinear localized dissipations in bounded domains has been investigated in Tcheugoué Tébou [9], Martinez [4] etc. and we can expect some fruitful results for unbounded domains. However, concerning the decay property of the wave equation with nonlinear dissipations in the whole space \mathbb{R}^N very few results seem to be known. K. Mochizuki and T. Motai [5] treated the case $\rho(x, u_t) = |u_t|^r u_t, 0 < r \leq 2/(N-2)^+$, and proved a logarithmic decay estimate

$$E(t) \equiv \frac{1}{2} \int_{\mathbb{R}^N} (|u_t(t)|^2 + |\nabla u(t)|^2) dx \leq C \{\log(1+t)\}^{-\nu}, \nu > 0.$$

K. Ono [8] considered the case $\rho(x, u_t) = u_t + |u_t|^r u_t, 0$

$$E(t) \leq C(\|u_0\|_{H_2}, \|u_1\|_{H_1})(1+t)^{\nu-1}$$

with $\nu = \max\{0, (4 - Nr)/2(r+2)\}$. This is refined if $(u_0, u_1) \in H_2 \cap L^p \times H_1 \cap L^p$. To our knowledge, these are the only results for the whole space, though the Klein-Gordon type wave equation with nonlinear dissipations like $|u_t|^r u_t$ have been treated by Nakao [6] and also Mochizuki and Motai [5].

2. SOME IDENTITIES AND THE BASIC INEQUALITY

Throughout this work we shall use only familiar Sobolev spaces and the definition of them are omitted. Assume Hyp.A. Then, for each $(u_0, u_1) \in H_1^0(\Omega) \times L^2(\Omega)$ the problem (1.1)-(1.3) has a unique solution $u(t)$ in $C_1([0, \infty); L^2(\Omega)) \cap C([0, \infty); H_1^0(\Omega))$, and for each $(u_0, u_1) \in H_1^0(\Omega) \cap H^2(\Omega) \times H_1^0(\Omega)$ the solution belongs to $X_{loc}^2 \equiv W_{loc}^{2,\infty}([0, \infty); L^2(\Omega)) \cap W_{loc}^{1,\infty}([0, \infty); H_1^0(\Omega)) \cap L_{loc}^\infty([0, \infty); H_2(\Omega))$, satisfying the estimate

$$(2.1) \quad \|u_{tt}\| + \|\nabla u_t\| + \|\Delta u(t)\| \leq C(\|\nabla u_1\| + \|\Delta u_0\| + \|\rho(x, u_1)\|) \equiv K_0 < \infty$$

for $0 \leq t < \infty$ and some $C > 0$.

These results are standard and well known (cf. Lions and Strauss [3]).

Let us consider the wave equation with a forcing term :

$$(2.2) \quad u_{tt} - \Delta u + \rho(x, u_t) = f(x, t) \quad \text{in } \Omega \times [0, \infty)$$

$$(2.3) \quad u(x, 0) = u_0, \quad u_t(x, 0) = u_1 \quad \text{in } \Omega$$

$$(2.4) \quad u(x, t) = 0 \quad \text{on } \partial\Omega \times [0, \infty).$$

For a moment we assume that $f \in W_{loc}^{1,2}([0, \infty); L^2(\Omega))$ and $u(t)$ is a solution in X_{loc}^2 of the problem (2.2)-(2.4).

For such solutions the following familiar identities hold. (See Lions[2] and Komornik[1]).

Lemma 2.1. *Let $h(x) = (h_1(x), h_2(x), \dots, h_N(x)) \in (W^{1,\infty}(\Omega))^N$ and $\phi(x) \in W^{1,\infty}(\Omega)$. Then,*

$$(2.5) \quad \frac{d}{dt} E(t) + \int_{\Omega} \rho(x, u_t) u_t dx = \int_{\Omega} f u_t dx,$$

$$\frac{d}{dt} \int_{\Omega} \varphi(x) u_t u dx - \int_{\Omega} \varphi(x) |u_t|^2 dx + \int_{\Omega} \nabla u \cdot \nabla(\varphi u) dx$$

$$+ \int_{\Omega} \varphi(x) \rho(x, u_t) u dx,$$

$$(2.6) \quad = \int_{\Omega} f \varphi(x) u dx$$

and

$$(2.7) \quad \frac{d}{dt} \left\{ \int_{\Omega} u_t h(x) \cdot \nabla u dx \right\} + \frac{1}{2} \int_{\Omega} \nabla \cdot h(x) |u_t|^2 dx$$

$$- \frac{1}{2} \int_{\Omega} \nabla \cdot h(x) |\nabla u|^2 dx + \int_{\Omega} \sum_{i,j=1}^N \frac{\partial h_i}{\partial x_j} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} dx$$

$$- \frac{1}{2} \int_{\Gamma(x_0)} \left| \frac{\partial u}{\partial \nu} \right|^2 \nu \cdot h(x) d\Gamma + \int_{\Omega} \rho(x, u_t) h(x) \cdot \nabla u dx$$

$$= \int_{\Omega} f h(x) \cdot \nabla u dx.$$

We denote by C positive generic constants independent of the initial data and f . $C(s)$ will denote a function of s which depends on s significantly. Our basic inequality reads as follows:

Proposition 2.2. *There exists $T_0 > 0$ such that if $T > T_0$ and $k > 0$ is sufficiently large, then*

$$(2.8) \quad X(t+T) - X(t) + k \int_t^{t+T} \int_{\Omega} \rho(x, u_t) u_t dx ds + \varepsilon_1 \int_t^{t+T} E(s) ds$$

$$\leq C \int_t^{t+T} \left\{ \int_{\omega} |u_t|^2 dx + \int_{\Omega_L} |\rho(x, u_t)|^2 dx + \int_{\Omega} (|f||u| + |f|^2) dx \right\} ds$$

with some $C > 0$, and $\varepsilon_1 > 0$, where $X(t)$ is a certain quantity equivalent to $E(t) + \|u(t)\|^2$.

The first term of the right-hand side is the most delicate one to control. But, we can show the following

Lemma 2.3. *There exists $T_0 > 0$, independent of u , such that if $T > T_0$, then the inequality*

$$(2.9) \quad \int_t^{t+T} \int_{\Omega_L} |u|^2 dx dt \leq C_\varepsilon(T) \int_t^{t+T} \left\{ \int_{\Omega} (|\rho(x, u_t)|^2 + |f|^2) dx + \int_{B_L^c \cup \omega} |u_t|^2 dx \right\} ds \\ + \varepsilon \int_t^{t+T} E(s) ds$$

holds for any fixed number $\varepsilon > 0$.

Now, we set

$$(2.10) \quad X(t) = \int_{\Omega} u_t(\phi(r)(x - x_0) - C_0 h(x)) \cdot \nabla u dx + \int_{\Omega} (\alpha + \eta^2) u_t u dx \\ + \frac{1}{2} \int_{\Omega_R^c} \alpha a(x) |u|^2 dx + kE(t).$$

It remains to show that $X(t)$ is equivalent to $E(t) + \|u(t)\|^2$. Indeed, we have :

Lemma 2.4. *For a large $k > 0$, there exist constants $C_1 > 0$ and $C_2 > 0$ such that for any $t \geq 0$,*

$$(2.11) \quad C_1(E(t) + \|u(t)\|^2) \leq X(t) \leq C_2(E(t) + \|u(t)\|^2).$$

3. MAIN RESULTS

We shall prove the bound of $\|u(t)\|$ and decay estimate of the energy $E(t) \equiv \frac{1}{2}(\|u_t(t)\|^2 + \|\nabla u(t)\|^2)$ as follows:

Theorem. *Let $u(t) \in X_{loc}^2$ be a solution of (1.1)-(1.3) satisfying (2.1). Under Hyp.A we have the estimates*

$$\|u(t)\|^2 \leq C(Q_0 + \|u_0\|^2)(t+1)^\nu \text{ and } E(t) \leq C(Q_0 + \|u_0\|^2)(t+1)^{\nu-1},$$

where ν and Q_0 are defined corresponding to the cases as follows:

(1) *If $0 \leq p < \infty$ and $0 \leq q \leq 2/(N-2)^+$ ($0 \leq q < \infty$ if $N = 1, 2$), then*

$$\nu = \max\{p/(p+2), q(N-2)^+/(4-(N-2)^+q)\},$$

$$(\nu = \max\{p/(p+2), \varepsilon\}, 0 < \varepsilon \ll 1, \text{ if } N = 2)$$

and

$$Q_0 = E(0)^{2/(p+2)} + E(0) + (K_0^2 + E(0))^{(q+1)\theta} E(0)^{2(q+1)(1-\theta)/(q+2)},$$

where $\theta = Nq/(q+1)(4-(N-2)^+q)$ ($\theta = (q+2)\varepsilon/(q+1)$ if $N = 2$).

(2) *If $-1 < p \leq 0$ and $0 \leq q \leq 2/(N-2)^+$, then*

$$\nu = \max\{-p/(p+2), q(N-2)^+/(4-(N-2)^+q)\},$$

$$(\nu = \max\{-p/(p+2), \varepsilon\}, 0 < \varepsilon \ll 1, \text{ if } N = 2)$$

and

$$Q_0 = E(0) + (K_0^2 + E(0))^{(q+1)\theta} E(0)^{2(q+1)(1-\theta)/(q+2)},$$

where $\theta = Nq/(q+1)(4-(N-2)^+q)$ ($\theta = (q+2)\varepsilon/(q+1)$ if $N = 2$).

(3) If $0 \leq p < \infty$ and $-1 < q \leq 0$, then

$$\nu = p/(p+2) \quad (\nu = \max\{p/(p+2), \varepsilon\} \text{ if } N = 2)$$

and

$$Q_0 = E(0)^{2/(p+2)} + E(0) + (K_0^2 + E(0))^{(q+1)\theta} E(0)^{4/(4-(N-2)^+q)}$$

with $\theta = -Nq/(q+1)(4-(N-2)^+q)$ ($\theta = \varepsilon/(q+1)$ if $N = 2$).

(4) If $-1 < p \leq 0$ and $-1 < q \leq 0$, then

$$\nu = \max\{-p/(p+2), q(N-2)^+/(4-(N-2)^+q)\},$$

$$(\nu = \max\{-p/(p+2), \varepsilon\}, 0 < \varepsilon \ll 1, \text{ if } N = 2)$$

and

$$Q_0 = E(0)^{2/(p+2)} + E(0) + (K_0^2 + E(0))^{(q+1)\theta} E(0)^{4/(4-(N-2)^+q)}$$

with $\theta = -Nq/(q+1)(4-(N-2)^+q)$ ($\theta = \varepsilon/(q+1)$ if $N = 2$).

The estimates are valid also for the Cauchy problem, $\Omega = \mathbb{R}^N$, if $N \geq 3$, where the boundary condition $u|_{\partial\Omega} = 0$ should be dropped.

REFERENCES

- [1] V. Komornik, Exact Controllability and Stabilization-The Multiplier Method, John Wiley, New York/Masson, Paris, 1994.
- [2] J. L. Lions, Contrôlabilité exacte, perturbations et stabilisation de systèmes distribués. Tome 1. Contrôlabilité exacte, RMA 8, Masson, 1988.
- [3] J. L. Lions and W. A. Strauss, Some non-linear evolution equations, *Bull. Soc. Math. France* **93**(1965), 43–96.
- [4] P. Martinez, A new method to obtain decay rate estimates for dissipative systems with localized damping, *Rev. Mat. Comp.* **12**(1999), 251-283.
- [5] K. Mochizuki and T. Motai, On energy decay-nondecay problems for the wave equations with nonlinear dissipative term in R^N , *J. Math. Soc. Japan* **47**(1995), 405-421.
- [6] M. Nakao, Energy decay of the wave equation with a nonlinear dissipative term, *Funkcial. Ekvac.* **26**(1983), 237–250.
- [7] M. Nakao, Energy decay for the linear and semilinear wave equations in exterior domains with some localized dissipations, *Math. Z.* **238**(2001), 781-797.
- [8] K. Ono, The time decay to the Cauchy problem for semilinear dissipative wave equations, *Adv. Math. Sci. Appl.* **9**(1999), 243–262.
- [9] L. R. Tcheugoué Tébou, Stabilization of the wave equation with localized nonlinear damping, *J. Differential Equations* **145**(1998), 502-524.

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