

## FINITE ADDITIVE MEASURE AND HYPERBOLIC SPACE

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ABSTRACT. We can extend the hyperbolic space beyond the infinity of the space by finite additive measure theory. This extended space which contains hyperbolic space as a subset has many natural and surprising properties which are similar to the case of spherical geometry.

### 0. Introduction

There are many hyperbolic models representing hyperbolic space  $\mathbb{H}^n$ , for example, upper half space model, unit disk model, hyperboloid model and Kleinian model. The Kleinian model  $K^n$  which is our matter of concern has the straight lines as geodesics. Here the Kleinian model is only defined at the inside of an unit ball by a well known Riemannian metric

$$ds^2 = \left( \frac{\sum x_i dx_i}{1 - |x|^2} \right)^2 + \frac{\sum dx_i^2}{1 - |x|^2}.$$

This Riemannian metric induces a volume form

$$dV = \frac{dx_1 \wedge \cdots \wedge dx_n}{(1 - |x|^2)^{\frac{n+1}{2}}}.$$

Here we can reform the Kleinian metric to a Kleinian  $\epsilon$ -metric

$$ds_\epsilon^2 = \left( \frac{\sum x_i dx_i}{d_\epsilon^2 - |x|^2} \right)^2 + \frac{\sum dx_i^2}{d_\epsilon^2 - |x|^2},$$

where  $d_\epsilon = 1 - \epsilon i$ , then the Kleinian volume form is changed to a Kleinian  $\epsilon$ -volume form

$$dV_\epsilon = \frac{d_\epsilon dx_1 \wedge \cdots \wedge dx_n}{(d_\epsilon^2 - |x|^2)^{\frac{n+1}{2}}}.$$

In the Kleinian model, the volume is calculated by

$$\mu(U) = \int_U \frac{dx_1 \wedge \cdots \wedge dx_n}{(1 - |x|^2)^{\frac{n+1}{2}}}, \text{ for a set } U \text{ in } K^n.$$

By the way the Kleinian  $\epsilon$ -volume form induces a volume for any set  $U$  in  $\mathbb{R}^n$  which contains  $K^n$  as a subset by

$$(1) \quad \bar{\mu}(U) = \lim_{\epsilon \rightarrow 0} \int_U \frac{d_\epsilon dx_1 \wedge \cdots \wedge dx_n}{(d_\epsilon^2 - |x|^2)^{\frac{n+1}{2}}},$$

and the above two volume integrals coincide at the case that a set  $U$  is contained in  $K^n$ . So the second integral volume formula (1) is the generalization of the first integral formula. Particularly for a set which intersects transversally to the ideal boundary of the hyperbolic space, the volume of the set has a finite complex value.

The volumes on the extended hyperbolic space have the hyperbolic invariant property as well as the hyperbolic space, i.e.,  $\overline{\mu}(f(U)) = \overline{\mu}(U)$  for an arbitrary hyperbolic isometry  $f$ . However the new measure  $\overline{\mu}$  does not admit the countable additive property. So we need find a suitable measure theory for this case: finite additive measure theory with an algebra (change countable union property to finite one) and a finite additive measure. Though the finite additive measure is weaker than an usual measure, we can construct a sufficient and natural theory on the extended Kleinian model  $\overline{K}^n$ .

The distance between two points on the extended space was firstly defined by Schlenker [4] by cross ratio. There he considered only distance. But our  $\epsilon$ -technique can be adapted for any other geometric objects, for example, distance between two points (easily turned out to be same to the definition of Schlenker), lengths of a curves, angles,  $k$ -dimensional volumes and so on (see [1]).

### 1. Finite additive invariant measure on the extended hyperbolic space

In order to speak about finite additive measure theory on the extended space, we have to construct  $(\overline{K}^n, \mathcal{M}, \mu)$ , where  $\mathcal{M}$  is an algebra in  $\overline{K}^n$  and  $\mu$  is a finite additive measure defined on  $\mathcal{M}$ .

The finite additive measure  $\mu$  is already chosen as  $\overline{\mu}$  in (1). Hence the remaining parts are a construction of  $\mathcal{M}$  and the finite additivity of  $\overline{\mu}$  on  $\mathcal{M}$ .

Before the construction, let's define the four Borel subcollections  $\mathbf{U}_1, \mathbf{U}_2, \mathbf{U}_3$ , and  $\mathbf{U}_4$  on  $\overline{K}^n$  and represent an element of  $\mathbf{U}_i$  as  $U_i$ .  $\mathbf{U}_1$  and  $\mathbf{U}_2$  is a collection of Borel sets in  $K^n$  and  $\overline{K}^n \setminus K^n$ , respectively, with finite volume,  $\mathbf{U}_3$  is a collection of center cones whose vertices are origin with Borel set sections, and  $\mathbf{U}_4$  is a collection of  $U_3 \setminus (U_1 \cup U_2)$ .

And let  $\mathcal{M}$  be a smallest algebra generated by  $\mathbf{U}_1, \mathbf{U}_2, \mathbf{U}_3$  types, and  $\mathcal{M}'$  be a set of  $U_1 \cup U_2 \cup U_4$ , and  $\mathcal{M}''$  be a set of  $U_1 \dot{\cup} U_2 \dot{\cup} U_4$ . Then we have  $\mathcal{M} = \mathcal{M}' = \mathcal{M}''$  (see [2]), so we can imagine the geometric figures for members of  $\mathcal{M}$ . The algebra  $\mathcal{M}$  contains all Borel sets which intersect transversally to the ideal boundary of the hyperbolic space, and all members of  $\mathcal{M}$  have a finite volume (see [2]). The proof of finite additivity of  $\overline{\mu}$  on  $\mathcal{M}$  is also easy.

For disjoint  $o_i \in \mathcal{M}$ , we get

$$\begin{aligned} \overline{\mu}(o_1 \cup o_2) &= \lim_{\epsilon \rightarrow 0} \int_{o_1 \cup o_2} dV_\epsilon = \lim_{\epsilon \rightarrow 0} \left( \int_{o_1} dV_\epsilon + \int_{o_2} dV_\epsilon \right) \\ &= \lim_{\epsilon \rightarrow 0} \int_{o_1} dV_\epsilon + \lim_{\epsilon \rightarrow 0} \int_{o_2} dV_\epsilon = \overline{\mu}(o_1) + \overline{\mu}(o_2). \end{aligned}$$

For a center cone  $U_3$  with a radius  $R$  and a section  $B$ , the volume of  $U_3$  is

obtained by a different and useful method (see [2])

$$\begin{aligned}\bar{\mu}(U_3) &= \lim_{\epsilon \rightarrow 0} \int_{U_3} \frac{d_\epsilon dx_1 \wedge \cdots \wedge dx_n}{(d_\epsilon^2 - |x|^2)^{\frac{n+1}{2}}} \\ &= \lim_{\epsilon \rightarrow 0} \int_B \int_0^R \frac{d_\epsilon r^{n-1}}{(d_\epsilon^2 - r^2)^{\frac{n+1}{2}}} dr d\theta \\ &= \int_B \int_\gamma \frac{r^{n-1}}{(1 - r^2)^{\frac{n+1}{2}}} dr d\theta,\end{aligned}$$

where  $d\theta$  is the volume form of the Euclidean unit sphere  $\mathbb{S}^{n-1}$  and  $\gamma$  represents a contour from 0 to  $R$  in  $\mathbb{C}$  traversing in a clockwise direction around 1, i.e.,  $\gamma$  is an over-going path with respect to 1.

A hyperbolic sphere model  $\mathbb{S}_H^n$  which is topologically equivalent to an unit sphere and considered a two fold covering of  $\overline{K}^n$  is better than the extended Kleinian model  $\overline{K}^n$ , when we consider invariance. For any transformation  $f$  in Lorentz group  $O(n, 1)$  which contains the hyperbolic isometry group as an index two subgroup and any set  $U$  in  $\mathcal{M}$ , we obtain  $\bar{\mu}(f(U)) = \bar{\mu}(U)$ . This invariance is an important property and difficult to prove (see [2]).

## 2. Extended hyperbolic space

In this section, we study the hyperbolic sphere model  $\mathbb{S}_H^n$  and the extended Kleinian model  $\overline{K}^n$  as an analytic continuation of the hyperbolic space  $\mathbb{H}^n$ . The hyperbolic sphere  $\mathbb{S}_H^n$  is topologically a sphere  $S^n$  which is obtained by projectivizing  $\mathbb{R}^{n,1}$  with only  $\mathbb{R}^+$ , the positive real numbers, we identifying a half lay emanating from the origin to a point. Hence  $\mathbb{S}_H^n$  is a double cover of  $\mathbb{R}P^n$  and consists of three parts: the hyperbolic part  $H^n$ , the Lorentzian part  $S_1^n$  and the projectivization of the light cone.

The metric on  $\mathbb{S}_H^n$  is given by the induced metric on  $H^n$  for the hyperbolic part  $H^n$  and the negative of the induced metric on the Lorentzian part  $S_1^n$ . Sometime, it would be helpful if we visualize  $\mathbb{S}_H^n$  as the Euclidean sphere in  $\mathbb{R}^{n,1}$  but with the induced metric coming from the Minkowski  $(\pm)$ -unit sphere through the radial projection and two fold covering of  $\overline{K}^n$ .

Note that our metric is negative of the original Lorentzian metric and the norm square of a tangent vector is positive. Hence this computation suggests us to choose a negative sign for a natural choice of sign for the norm of a tangent vector in Lorentzian part of  $\mathbb{S}_H^n$ , and for tangent vectors in the radial direction in Lorentzian part of general  $\mathbb{S}_H^n$  as well. Similarly it is not hard to check the clockwise contour integral of the volume form gives sign  $-i^{n-1}$  for Lorentzian part. If  $n = 2$ , the sign for the volume is  $-i$  and this suggests us that  $i$  is the natural choice for the sign for the norm of the spacelike tangent vector in the Lorentzian part of  $\mathbb{S}_H^2$  and in  $\mathbb{S}_H^n$  as well. This gives us a natural choice of signs for the norm of various tangent vectors which are compatible with volume form on  $\mathbb{S}_H^n$ . We summarize as follows.

A tangent vector on the hyperbolic part on  $\mathbb{S}_H^n$  has a positive real norm, and a tangent vector on the Lorentzian part on  $\mathbb{S}_H^n$  has a negative real, zero, or positive pure imaginary norm depending on whether it is timelike, null, or spacelike respectively.

We can see various similarities between the Euclidean sphere  $\mathbb{S}^n$  and the hyperbolic sphere  $\mathbb{S}_H^n$  in the following listed theorems (see [1], [3]).

**Theorem 1.**  $\text{vol}(\mathbb{S}_H^n) = i^n \cdot \text{vol}(\mathbb{S}^n)$ .

*Remark.* If we change the  $d$  as  $1 + \epsilon i$ , we have different relation between  $\text{vol}(\mathbb{S}_H^n)$  and  $\text{vol}(\mathbb{S}^n)$ :  $\text{vol}(\mathbb{S}_H^n) = (-i)^n \text{vol}(\mathbb{S}^n)$ .

**Theorem 2.** *The total length of any great circle is  $2\pi i$ .*

**Theorem 3.** *The  $k$ -dimensional volume of any  $k$ -dimensional geodesic sphere is  $\text{vol}(\mathbb{S}_H^k)$ .*

**Theorem 4.** *For any vectors  $x$  and  $y$  in  $\mathbb{S}_H^n$  and a point  $p$  in  $x^\perp \cap y^\perp$ , we have*

$$\begin{aligned}\langle x, y \rangle &= \|x\| \|y\| \cos \angle(x_p, y_p) \\ \langle x, y \rangle &= \|x\| \|y\| \cosh d_H(x, y).\end{aligned}$$

**Theorem 5.** *The area of a triangle in  $\mathbb{S}_H^2$  is  $\pi - A - B - C$ , where  $A, B, C$  are the angles of the given triangle.*

**Theorem 6.** *Letting  $A, B, C$  stand for the angles and  $a, b, c$  for the extended hyperbolic length of opposite sides, we obtain the hyperbolic cosine law and the dual cosine law on the hyperbolic sphere  $\mathbb{S}_H^2$ ,*

$$\begin{aligned}\cos C &= \frac{\cosh a \cosh b - \cosh c}{\sinh a \sinh b} \\ \cosh c &= \frac{\cos A \cos B + \cos C}{\sin A \sin B}.\end{aligned}$$

*Also we have the spherical cosine law and dual cosine law on the spherical sphere  $\mathbb{S}_S^2$ , where  $a, b, c$  represent the extended spherical length,*

$$\begin{aligned}\cos C &= \frac{\cos c - \cosh a \cosh b}{\sin a \sin b} \\ \cos c &= \frac{\cos A \cos B + \cos C}{\sin A \sin B}.\end{aligned}$$

**Theorem 7.** *Letting  $A, B, C$  stand for the angles and  $a, b, c$  for the extended hyperbolic length of opposite sides, we obtain the hyperbolic sine law on the hyperbolic sphere  $\mathbb{S}_H^2$ ,*

$$\frac{\sinh a}{\sin A} = \frac{\sinh b}{\sin B} = \frac{\sinh c}{\sin C}.$$

*Also we have the spherical sine law on the spherical sphere  $\mathbb{S}_S^2$ , where  $a, b, c$  represent the extended spherical length,*

$$\frac{\sin a}{\sin A} = \frac{\sin b}{\sin B} = \frac{\sin c}{\sin C}.$$

*Remark.* The spherical sphere  $\mathbb{S}_S^2$  is made by the following  $\epsilon$ -metric on  $\mathbb{R}^n$ :

$$ds_\epsilon^2 = - \left( \frac{\sum x_i dx_i}{d_\epsilon^2 - |x|^2} \right)^2 - \frac{\sum dx_i^2}{d_\epsilon^2 - |x|^2}.$$

## REFERENCES

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